## Defining the Orientation of the ${ }^{15} \mathrm{~N}$ Shielding Tensor Using ${ }^{15} \mathrm{~N}$ NMR Relaxation Data for a Protein in Solution

Jonathan Boyd and Christina Redfield*<br>Department of Biochemistry and Oxford Centre for Molecular Sciences, Oxford University South Parks Road, Oxford, OX1 3QT, U.K.<br>Received May 6, 1998<br>Revised Manuscript Received July 20, 1998

${ }^{15} \mathrm{~N}$ relaxation studies provide information about protein backbone dynamics in solution. ${ }^{1-3}$ These data are generally interpreted using the "model-free" formalism ${ }^{4}$ which yields a generalized order parameter, $S^{2}$, and a time constant for internal motion, $\tau_{\mathrm{e}}$. Analysis of these values allows the distinction between backbone mobility on a picosecond to nanosecond time scale and static disorder in NMR-derived structures. Recently, it has been demonstrated that ${ }^{15} \mathrm{~N}$ relaxation data are sensitive to rotational diffusion anisotropy;, ${ }^{3,5-9}$ variations in experimental $T_{1} /$ $T_{2}$ ratios can be used to define the orientation and magnitude of the diffusion tensor. As a result, ${ }^{15} \mathrm{~N} T_{1} / T_{2}$ ratios can be used as constraints for NMR structure refinement ${ }^{8}$ or to provide information about the conformation of domains in multidomain proteins. ${ }^{5}$

Two methods ${ }^{5-9}$ have emerged for defining the orientation of the rotational diffusion tensor in the molecular frame using experimental ${ }^{15} \mathrm{~N}$ relaxation data. In the first method, ${ }^{5,9}$ a local diffusion coefficient, $D_{i}$, is determined for each residue $i$. In the second, ${ }^{6-8}$ a direct fitting to the experimental ${ }^{15} \mathrm{~N} T_{1} / T_{2}$ ratios is performed. $T_{1} / T_{2}$ ratios are calculated for spherical, symmetric, and asymmetric diffusion models, and the goodness-of-fit to experimental data is assessed using the $\chi^{2}$ statistic:

$$
\begin{equation*}
\chi^{2}=\sum_{i=1}^{n}\left[\left(T_{1 i} / T_{2 i}\right)^{\text {calcd }}-\left(T_{1 i} / T_{2 i}\right)^{\text {expt } 2}\right]^{2} / \Delta^{2} \tag{1}
\end{equation*}
$$

where $\left(T_{1 i} / T_{2 i}\right)^{\text {calcd }}$ and $\left(T_{1 i} / T_{2 i}\right)^{\text {expt }}$ are the calculated and experimental values, respectively, for residue $i, \Delta$ is the experimental uncertainty in $T_{1} / T_{2}$, and $n$ is the number of residues. The selection of a spherical, symmetric, or asymmetric diffusion model for a protein, where each model has an increasing number of fitted parameters, is made by assessing the statistical significance of the additional parameters using the $F$-test. ${ }^{6,9,10}$

Solid-state NMR studies ${ }^{11}$ and quantum mechanical calculations ${ }^{12}$ of peptides and peptide analogues have shown that the ${ }^{15} \mathrm{~N}$ shielding tensor is nearly axially symmetric with its principal

[^0]axis located in the peptide plane. These studies indicate that the principal axis is not collinear with the ${ }^{15} \mathrm{~N}-\mathrm{H}$ bond vector but is inclined away from the ${ }^{15} \mathrm{~N}-\mathrm{H}$ bond toward the ${ }^{15} \mathrm{~N}-\mathrm{C}^{\prime}$ bond by an angle $\alpha$. A number of different values for $\alpha$, typically ranging from 15 to $25^{\circ}$, have been reported from these studies. The value of $\alpha$ has also been investigated in a recent solution NMR study. ${ }^{13}$ The small magnetic-field-dependent chemical-shift differences ${ }^{13}$ observed from partially oriented anisotropic molecules depend on $\alpha$; an analysis of these shift differences in a zinc-finger-DNA complex has found $\alpha$ to be in the range $13 \pm 5^{\circ}$. ${ }^{15} \mathrm{~N}$ chemical shift anisotropy (CSA) and ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ dipolar cross-correlation is influenced by the value of $\alpha .{ }^{14}$ A detailed analysis of the CSA dipolar cross-correlation in ubiquitin has been presented; ${ }^{14}$ the angle $\alpha$ was taken to be $20^{\circ}$ for all ${ }^{15} \mathrm{~N}$ nuclei.

The influence of $\alpha$ on the calculated ${ }^{15} \mathrm{~N} T_{1}$ and $T_{2}$ relaxation times, when an anisotropic rotational diffusion model is employed, has not previously been considered. Hitherto, calculations of these relaxation times, involving autocorrelation spectral densities, ${ }^{15,16}$ have employed a single set of angles to define the orientation of both the ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ internuclear vector and the principal component of the ${ }^{15} \mathrm{~N}$ shielding tensor with respect to the principal axes of the diffusion tensor, implicitly assuming that $\alpha$ is zero. Here, we show that the orientation of the ${ }^{15} \mathrm{~N}$ shielding tensor, as defined by $\alpha$, has an influence upon calculated $T_{1} / T_{2}$ ratios and that this effect is detectable in experimental data.
The ratio $\left(T_{1 i} / T_{2 i}\right)$ calcd for the $i$ th ${ }^{15} \mathrm{~N}$ nucleus is calculated using two sets of spectral density functions defined by:

$$
\begin{equation*}
J_{i k}(\omega)=\sum_{j=1}^{n} A_{i j k}\left[\tau_{j} /\left(1+\omega^{2} \tau_{j}^{2}\right)\right] \tag{2}
\end{equation*}
$$

where $n=1,3$, and 5 for the spherical, symmetric, and asymmetric top models, respectively. The spectral densities $J_{i 1^{-}}$ $(\omega)$ are used for ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ dipolar contributions, and $J_{i 2}(\omega)$ for ${ }^{15} \mathrm{~N}$ CSA contributions to the $T_{1}$ and $T_{2}$ relaxation equations, ${ }^{15,16}$ respectively. The time constants $\tau_{j}$ depend on the principle values of the diffusion tensor, as defined previously. ${ }^{15,16}$ The coefficients $A_{i j 1}$, for the dipolar terms, depend on the orientation of the ${ }^{15} N_{(i)}-$ ${ }^{1} \mathrm{H}_{(i)}$ internuclear vector with respect to the diffusion tensor whereas the coefficients $A_{i j 2}$, for the CSA terms, depend on the orientation of the ${ }^{15} \mathrm{~N}_{(i)}$ shielding tensor with respect to the diffusion tensor. ${ }^{15,16}$ In the case of a symmetric top $A_{i j 1}$ depends on the angle $\theta$ between the ${ }^{15} \mathrm{~N}_{(i)}-{ }^{-1} \mathrm{H}_{(i)}$ vector and the principal axis of the diffusion tensor while $A_{i j 2}$ depends on the angle $\theta^{\prime}$ between the principal axis of the ${ }^{15} \mathrm{~N}$ shielding tensor and the principal
(11) (a) Munowitz, M.; Aue, A. P.; Griffin, R. G. J. Chem. Phys. 1982, 77, 1686-1689. (b) Harbison, G. S.; Jelinski, L. W.; Stark, R. E.; Torchia, D. A.; Herzfeld, J.; Griffin, R. G. J. Magn. Reson. 1984, 60, 79-82. (c) Hartzell, C. J.; Pratum, T. K.; Drobny, G. P. J. Chem. Phys. 1987, 87, 43244331. (d) Oas, T. G.; Hartzell, C. J.; Dahlquist, F. W.; Drobny, G. P. J. Am. Chem. Soc. 1987, 109, 5962-5966. (e) Hiyama, Y.; Niu, C.-H.; Silverton, J. V.; Bavoso, A.; Torchia, D. A. J. Am. Chem. Soc. 1988, 110, 2378-2383. (f) Teng, Q.; Cross, T. A. J. Magn. Reson. 1989, 85, 439-447. (g) Shoji, A.; Ozaki, T.; Fujito, T.; Deguchi, K.; Ando, S.; Ando, I. Macromolecules 1989, 22, 2860-2863. (h) Mai, W.; Hu, W.; Wang, C.; Cross, T. A. Protein Sci. 1993, 2, 532-542. (i) Wu, C. H.; Ramamoorthy, A.; Gierasch, L. M.; Opella, S. J. J. Am. Chem. Soc. 1995, 117, 6148-6149.
(12) (a) Facelli, J. C.; Pugmire, R. J.; Grant, D. M. J. Am. Chem. Soc. 1996, 118, 5488-5489. (b) Walling, A. E.; Pargas, R. E.; de Dios, A. C. J. Phys. Chem. A 1997, 101, 7299-7303.
(13) Ottiger, M.; Tjandra, N.; Bax, A. J. Am. Chem. Soc. 1997, 119, 98259830.
(14) Tjandra, J.; Szabo, A.; Bax, A. J. Am. Chem. Soc. 1996, 118, 69866991.
(15) (a) Woessner D. T. J. Chem. Phys. 1962, 252, 7740-7743. (b) Werbelow, L. G.; Grant, D. M. In Advances in Magnetic Resonance; Waugh, J. S., Ed.; Academic Press: New York, 1977; Vol. 9, pp 189-299.
(16) Spiess, H. W. In NMR Basic Principles and Progress; Diehl P., Fluck, E., Kosfeld, R.. Eds.; Springer-Verlag: New York, 1978; Vol. 15, pp 55214.

Table 1. Comparison of $\chi^{2}$ Values Obtained Using a Symmetric Diffusion Model with the ${ }^{15} \mathrm{~N}$ Shielding Tensor Angle $\alpha$ Set to $0^{\circ}$ (left) and with the Angle $\alpha$ Fitted (right) for Four Lysozyme Crystal Structures

| structure $^{a}$ | $\chi^{2 b}$ | $D_{\\|} / D_{\perp}$ | $\alpha(\mathrm{deg})$ | $\chi^{2 c}$ | $D_{\\|} / D_{\perp}$ | $\alpha(\mathrm{deg})$ | $F-$ value ${ }^{d}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5LYM | 140.2 | $1.29 \pm 0.01$ | 0.0 | 124.1 | $1.32 \pm 0.02$ | $28.5 \pm 7.6$ |  |
| 1UCO | 142.8 | $1.29 \pm 0.01$ | 0.0 | 128.1 | $1.32 \pm 0.02$ | $27.2 \pm 8.0$ |  |
| 3LYM | 145.2 | $1.28 \pm 0.01$ | 0.0 | 123.7 | $1.32 \pm 0.02$ | $32.2 \pm 7.1$ |  |
| 6LYT | 146.6 | $1.29 \pm 0.01$ | 0.0 | 126.7 | $1.32 \pm 0.02$ | $33.9 \pm 8.4$ |  |

${ }^{a}$ X-ray coordinates for hen lysozyme ${ }^{19}$ were obtained from the Brookhaven Protein Databank. Amide protons $\left(\mathrm{H}^{\mathrm{N}}\right)$ were added to the X-ray structures using X-PLOR. ${ }^{20}$ The structure was then energy minimized with the backbone heavy atoms constrained to remain at their initial positions. This energy minimization was necessary to ensure that $\mathrm{H}^{\mathrm{N}}$ was located in the peptide plane. A vector representing the principal component of the ${ }^{15} \mathrm{~N}$ shielding tensor for each residue $i+1$ was generated using the coordinates of the atoms $\mathrm{H}^{\mathrm{N}}{ }_{(i+1)}, \mathrm{N}_{(i+1)}$, and $\mathrm{C}^{\prime}(i)$, this vector lies in the plane defined by these atoms, inclined at an angle $\alpha$ to the ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ bond. $\mathrm{A}{ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ internuclear distance of 0.102 nm was used in calculations of ${ }^{15} \mathrm{~N}$ $T_{1}$ and $T_{2}$ relaxations times. ${ }^{b}$ The angle $\alpha$ is fixed at $0^{\circ}$ and the parameters $D_{\| \|}$and $D_{\perp}$ and the orientation of the diffusion tensor are fitted to minimize $\chi^{2}$. A total of $121 T_{1} / T_{2}$ ratios have been measured experimentally at 50.6 MHz for lysozyme. ${ }^{2}$ Residues with significant motions on the picosecond time scale, exchange broadening, or found in certain surface loops involved in crystal contacts have been removed from analysis; 6,18 93 $T_{1} / T_{2}$ ratios are used in the present study. ${ }^{c} D_{\|}, D_{\perp}$, the orientation of the diffusion tensor, and $\alpha$ are fitted to minimize $\chi^{2}$. The error in $D_{\|} / D_{\perp}$ and $\alpha$ has been estimated from 500 Monte Carlo simulations. The effective correlation time, $\tau_{\mathrm{r}}=1 /\left(2 D_{\|}+4 D_{\perp}\right)$, decreases by at most 0.002 ns when $\alpha$ is fitted. A value of 5.78 ns is found for all structures. ${ }^{d}$ The $F$-test indicates that the symmetric diffusion model with the additional parameter $\alpha$ is a statistically significant improvement over the symmetric diffusion model with $\alpha=0^{\circ}$.


Figure 1. Plot of $\chi^{2}$ as a function of the ${ }^{15} \mathrm{~N}$ shielding tensor angle $\alpha$. For each value of $\alpha$, the parameters $D_{\|}, D_{\perp}$ and the orientation of the diffusion tensor were fitted. Calculations using uniform ${ }^{15} \mathrm{~N}$ CSA values of either $-160,-170$, or -180 ppm are shown by short-dashed, solid, and long-dashed lines, respectively. The vertical dashed lines indicate the range of errors in the fitted value of $\alpha$ found using Monte Carlo simulations for a ${ }^{15} \mathrm{~N}$ CSA value of -170 ppm . The inset shows the angles $\alpha, \theta$, and $\theta^{\prime}$ which define the relative orientations of the principal component of the ${ }^{15} \mathrm{~N}$ shielding tensor $\left(\sigma_{\|}\right)$, the ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ internuclear vector, and the principal component of the diffusion tensor $\left(D_{\|}\right)$.
axis of the diffusion tensor. The value of $\theta^{\prime}$ falls in the range of $\theta \pm \alpha$, the exact value depending on the orientation of the peptide plane with respect to the diffusion tensor (Figure 1). For the calculations performed here, we have assumed that the ${ }^{15} \mathrm{~N}$ chemical shift asymmetry parameter is zero, that any contribution to ${ }^{15} \mathrm{~N}$ relaxation from the antisymmetric component of the shielding tensor ${ }^{16,17}$ is small and can be neglected, and that the ${ }^{15} \mathrm{~N}$ shielding tensor experiences identical dynamic behavior to the ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ dipolar interaction.

The ${ }^{15} \mathrm{~N} T_{1}$ and $T_{2}$ relaxation data were obtained at 50.6 MHz from a uniformly ${ }^{15} \mathrm{~N}$-labeled sample of hen lysozyme as described previously. ${ }^{2}$ A recent analysis has shown that a symmetric rotational diffusion model best describes these data. ${ }^{18}$ The change in the $\chi^{2}$ statistic as the angle $\alpha$ is varied from 0 to $60^{\circ}$ for the 5 LYM lysozyme structure ${ }^{19}$ is shown in Figure 1. The agreement between the experimental and calculated $T_{1} / T_{2}$ ratios improves as $\alpha$ is increased, with a minimum at $28.5^{\circ}$. Similar fits have

[^1]been carried out using three other lysozyme X-ray structures, ${ }^{19}$ the results are summarized in Table 1. For these structures, the optimal value of $\alpha$ is found to vary between 27.2 and $33.9^{\circ}$. The $F$-values listed in Table 1 correspond to a confidence level of greater than $99.8 \%$ indicating that the model which includes the angle $\alpha$ gives a significant improvement over the more widely used model which assumes that $\alpha=0^{\circ}$. Monte Carlo simulations have been used to assess the error in $\alpha$ resulting from experimental uncertainty in the ${ }^{15} \mathrm{~N} T_{1}$ and $T_{2}$ values; errors of $\pm 7-8^{\circ}$ are found for each of the four structures. This fairly large error is to be expected because the variation of $\chi^{2}$ as a function of the angle $\alpha$ is found to exhibit a broad shallow minimum as shown in Figure 1. These values of $\alpha$ are obtained with a uniform ${ }^{15} \mathrm{~N}$ CSA value of $-170 \mathrm{ppm} .{ }^{14}$ If this value is changed to -160 or -180 ppm , $\alpha$ increases or decreases by $2.4^{\circ}$, respectively, as shown in Figure 1. These values for $\alpha$ are slightly larger than those obtained previously. ${ }^{11,13}$ The value of $D_{\|} / D_{\perp}$ is found to increase when the value of $\alpha$ is fitted. For example, $D_{\|} / D_{\perp}$ increases from 1.29 to 1.32 , for the 5 LYM structure, when $\alpha$ is increased from $0^{\circ}$ to $28.5^{\circ}$. If $D_{\| /} / D_{\perp}$ and the orientation of the diffusion tensor are fixed to the values found with $\alpha=0^{\circ}$ and only $\alpha$ is fitted, then somewhat lower values of $\alpha$ are found; $\alpha$ equals $21.9,21.0,23.9$, and $24.3^{\circ}$ for the $5 \mathrm{LYM}, 1 \mathrm{UCO}, 3 \mathrm{LYM}$, and 6LYT structures, respectively. These values for $\alpha$ are now closer to the values reported from solid-state $\mathrm{NMR}^{11}$ and from magnetic-fielddependent chemical-shift studies. ${ }^{13}$

Although the uncertainty in the value of $\alpha$ is quite large, the calculations for lysozyme show that improved agreement between the experimental and calculated $T_{1} / T_{2}$ ratios is obtained when using a symmetric rotational diffusion model in combination with the assumption that the ${ }^{15} \mathrm{~N}$ shielding tensor is not collinear with the ${ }^{15} \mathrm{~N}-{ }^{1} \mathrm{H}$ bond vector. The data analyzed here were collected at relatively low field for a protein of moderate anisotropy $\left(D_{\|} / D_{\perp}\right.$ $=1.3$ ). It is expected that the influence of $\alpha$ on $T_{1} / T_{2}$ ratios will be more pronounced at higher fields $(18.8 \mathrm{~T})$ and for proteins with a larger anisotropy. This has implications for the analysis of frequency-dependent ${ }^{15} \mathrm{~N}$ relaxation data and for the use of $T_{1} / T_{2}$ ratios in the refinement of protein structure.

Acknowledgment. This work is from the Oxford Centre for Molecular Sciences, supported by BBSRC, EPSRC and MRC. One of us (C.R.) thanks the BBSRC for an Advanced Research Fellowship.

Supporting Information Available: One table containing experimental and calculated $T_{1} / T_{2}$ ratios with $\alpha=0^{\circ}$ and $\alpha=28.5^{\circ}$ for the 5LYM lysozyme structure (4 pages, print/PDF). See any current masthead page for ordering information and Web access instructions.

## JA9815733

[^2]
[^0]:    (1) (a) Kay, L. E.; Torchia, D. A.; Bax, A. Biochemistry 1989, 28, 40004009. (b) Stone, M. J.; Fairbrother, W. J.; Palmer, A. G.; Reizer, J.; Saier, M. H.; Wright, P. E. Biochemistry 1992, 31, 4394-4406. (c) Redfield, C.; Boyd, J.; Smith, L. J.; Smith, R. A. G.; Dobson, C. M. Biochemistry 1992, 31, 10431-10437.
    (2) Buck, M.; Boyd, J.; Redfield, C.; MacKenzie, D. A.; Jeenes, D. J.; Archer, D. B.; Dobson, C. M. Biochemistry 1995, 34, 4041-4055.
    (3) Barbato, G.; Ikura, M.; Kay, L. E.; Pastor, R. W.; Bax, A. Biochemistry 1992, 31, 5269-5278.
    (4) Lipari, G.; Szabo, A. J. Am. Chem. Soc. 1982, 104, 4546-4570.
    (5) Bruschweiler, R.; Liao, X.; Wright, P. E. Science 1995, 268, 886889.
    (6) (a) Tjandra, N.; Feller, S. E.; Pastor, R. W.; Bax, A. J. Am. Chem. Soc. 1995, 117, 12562-12566. (b) Tjandra, N.; Wingfield, P.; Stahl, S.; Bax, A. J. Biomol. NMR 1996, 8, 273-284.
    (7) Phan, I. Q. H.; Boyd, J.; Campbell, I. D. J. Biomol. NMR 1996, 8 , 369-378.
    (8) Tjandra, N.; Garrett, D. S.; Gronenborn, A. M.; Bax, A.; Clore, G. M. Nat. Struct. Biol. 1997, 4, 443-449.
    (9) Lee, L. K.; Rance, M.; Chazin, W. J.; Palmer, A. G. J. Biomol. NMR 1997, 9, 287-298.
    (10) Bevington, P. R.; Robinson, D. K. Data Reduction and Error Analysis for the Physical Sciences, 2nd ed.; McGraw-Hill: New York, 1992; pp 194220.

[^1]:    (17) (a) Buckingham, A. D.; Malm, S. M. Mol. Phys. 1971, 22, 11271130. (b) Lynden-Bell, R. M. Mol. Phys. 1975, 29, 301-303. (c) Kowalewski, J.; Werbelow, L. J. Magn. Reson. 1997, 128, 144-148.
    (18) Redfield, C. Manuscript in preparation.
    (19) (a) Rao, S. T.; Sunaralingam, M. Acta Crystallogr. 1996, D52, 170175. (b) Nagendra, H. G.; Sudarsanakumar, C.; Vijayan, M. Acta Crystallogr. 1996, D 52, 1067-1074. (c) Kundrot, C. E.; Richards, F. M. J. Mol. Biol. 1987, 193, 157-170. (d) Young, A. C.; Tilton, R. F.; Dewan, J. C. J. Mol. Biol. 1994, 235, 302-317.

[^2]:    (20) Brunger, A. T. X-PLOR Version 3.1: A system for X-ray Crystallography and NMR; Yale University Press: New Haven, CT, 1992.

